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1	The dynamics of suspended particulate matter (SPM) and
2	chlorophyll-a from intratidal to annual time scales in a coastal
3	turbidity maximum
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# 23 Abstract

24 The analysis of 1.8 years of data gives an understanding of the response to varying forcing of suspended particulate matter (SPM) and chlorophyll-a (CHL-a) in a coastal turbidity maximum zone 25 26 (TMZ). Both temporal and vertical concentration variations in the near-bed layer (0 - 2 m) in the 27 shallow (11 m deep) coastal zone at 1 km off the Dutch coast are shown. Temporal variations in the 28 concentration of both parameters are found on tidal and seasonal scales, and a marked response to 29 episodic events (e.g. storms). The seasonal cycle in the near-bed CHL-a concentration is determined 30 by the spring bloom. The role of the wave climate as the primary forcing in the SPM seasonal cycle is 31 discussed. The tidal current provides a background signal, generated predominantly by local 32 resuspension and settling and a minor role is for advection in the cross-shore and the alongshore 33 direction. We tested the logarithmic Rouse profile to the vertical profiles of both the SPM and the 34 CHL-a data, with respectively 84% and only 2% success. The resulting large percentage of low Rouse 35 numbers for the SPM profiles suggest a mixed suspension is dominant in the TMZ, ie. surface SPM concentrations are in the same order of magnitude as near-bed concentrations. 36

37

# 38 Keywords

39 Wave climate // In-situ observations // Bed shear stress // SPM // chlorophyll-a // season

40

# 41 Highlights

42	•	The seasonal SPM variation is primarily driven by the seasonal wave climate
43	•	SPM and CHL-a concentrations have an opposed seasonal cycle

- The tidal current generates a background signal in both concentrations
- 45 The majority of SPM profiles fit to the Rouse profile

## 46 **1. Introduction**

47 Coastal areas all over the world are under constant pressure due to the growth of the world population and climate change. Currently implemented coastal adaptations to counteract the effects 48 of climate change may affect the local ecology and the surrounding ecosystem. In the Netherlands, 49 50 several engineering solutions have been developed and implemented to mitigate negative effects on 51 the ecosystem (Borsje et al., 2011; De Jong et al., 2015). Understanding the effect of these solutions 52 on the ecosystem requires basic knowledge of the coastal system in both the physical and the 53 ecological sense. One parameter to consider in the Dutch coastal zone is suspended particulate 54 matter (SPM), as the coastal zone forms a transport route of SPM from the Strait of Dover and rivers 55 towards the Wadden Sea. SPM is composed of organic and terrigenous particles and its 56 concentration and variation affects the pelagic as well as the benthic ecosystem.

57 In the southern North Sea, highest SPM concentrations are found close to the coasts as observed by many surface observations such as remote sensing data (Eleveld et al., 2008; Fettweis et 58 al., 2007; Van Raaphorst et al., 1998), and in-situ data (Dyer and Moffatt, 1998; Fettweis et al., 2007; 59 Howarth et al., 2002; Huthnance, 1997; Van Alphen, 1990; Visser et al., 1991). More specifically for 60 61 the Dutch coast, a near-bed turbidity maximum zone (TMZ) is present along the coastline at a 62 maximum distance of 3 km from the shore (Van Alphen, 1990; Van der Hout et al., 2015). Thus high 63 SPM concentrations are close to the coast and nearby engineering solutions. All along the coastline 64 SPM is accumulating in cross-shore direction in the TMZ (Visser et al., 1991) and is transported northwards (Van der Giessen et al., 1990) forming a "coastal river of silt". The majority of the 65 66 northward transport of SPM along the coast, and also of nutrients, plankton, fish larvae and other 67 juvenile marine fauna, may be concentrated in this small region. The TMZ is likely to control the distribution of macrobenthos, and the prey items of higher trophic levels such as predatory fish or 68 69 birds (Van der Veer et al., 2015), by its high turbidity and nutrient levels. It is in this same zone that 70 high densities of the benthic bivalve Ensis directus have been observed (Witbaard et al., 2015), which 71 is a filter feeder thriving on suspended (phytoplankton) particles. These bivalves may contribute to 72 the accumulation of SPM in the TMZ by (seasonal) burial of SPM into the seabed (Witbaard et al., 73 2016; this volume).

Processes that affect the levels of SPM concentrations in the southern North Sea are depth, tidal currents, biota, waves and density driven currents (Fettweis et al., 2014; Stanev et al., 2009; Visser et al., 1991). The tidal wave in this area propagates in an anti-clockwise direction and generates strong tidal currents near the coasts. In combination with waves, high bed shear stresses are generated (Stanev et al., 2009) and fine bottom particles are resuspended. Unraveling the response of SPM and CHL-a concentrations to the tidal regime, the waves or biological factors has 80 been subject of many field studies. Variations in SPM and CHL-a linked to tidal regimes were reported by Bass et al. (2002), Blauw et al. (2012), Joordens et al. (2001), McCandliss et al. (2002) and Orton 81 82 and Kineke (2001). Studies in Liverpool Bay and Jiaozhou Bay highlighted a wave-related dependence 83 of the SPM concentration (Fettweis et al., 2012; Krivtsov et al., 2012 and references therein). SPM 84 concentrations in the Dutch and Belgian coastal zone are higher in winter than in summer, which has 85 been linked to the seasonal pattern in storms (Eleveld et al., 2008; Fettweis et al., 2012; Visser et al., 1991). However, in recent research in the Belgian coastal zone (bordering the Dutch coast) this 86 87 seasonal pattern is hypothesized to be related to the seasonal cycle in the phytoplankton population 88 (Fettweis et al., 2014). The phytoplankton population is determined by light, nutrient availability, 89 temperature and grazing, and on the other hand the population affects flocculation of SPM (Jago et 90 al., 2007; Jones et al., 1998; Passow, 2002) and hence acts as a seasonal biological process affecting 91 the fall velocity of SPM.

92 The focus of this study is to describe the processes that control the temporal variation of 93 SPM and CHL-a concentrations in the near-bed water layer (0 - 2 m) in the TMZ along the Dutch 94 coast. The seasonal cycle is our main interest, but also smaller scales down to the tidal cycle will be 95 analyzed. The dataset for the analysis was obtained with an instrumented bottom frame deployed in 96 the TMZ for a period of almost 2 years. Details of the instrumentation and the specifics of the 97 location of bottom frame are given in Section 2, as well as details on the methods applied to analyze 98 the dataset. The time-series includes storm conditions, periods for which data is otherwise difficult to 99 obtain due to unworkable high swell or clouds obscuring satellite imagery. The effects of the wave 100 climate and the tide on the instant concentrations and on the seasonal cycle is presented in Section 101 3. Additionally, we will discuss the suspension mode over time on basis of the Rouse number,  $R_n$ . The 102 validity of the Rouse profile will first be discussed in section 3.3, as several of the assumptions for the 103 equilibrium profile are violated in the research area (Mehta, 2014; Orton and Kineke, 2001). If valid, 104 the near-bed measured concentrations can be extended up to the surface with the logarithmic Rouse 105 number (Mehta, 2014). We will use the Rouse number to estimate the suspension mode; either a 106 near-bed dominated suspension (high  $R_n$ ) or a mixed suspension (low  $R_n$ ) in which the SPM is more 107 uniformly distributed over the water column (Mehta, 2014). In Section 4 we will discuss the 108 (interaction of the) seasonal cycles of the SPM and CHL-a concentrations, the resuspension and 109 advective character of the tidal oscillation and the effect of the suspension mode for transport 110 estimations. In Section 5 we conclude with a description of the character of the SPM suspension near 111 the seabed in the coastal TMZ in which the tide and the wave climate each play a role.

## 112 **2. Methods**

# 113 **2.1.** Location and field observations

#### 114 2.1.1. Bottom frame location

The study area is located along the Dutch coast, 20 km north of the harbor of IJmuiden, in the 115 southern North Sea (Figure 1). In 2011 and 2012, a bottom frame was deployed 1.2 km offshore 116 117 (52°38.249'N; 4°36.294'E) for the project 'Monitoring and Evaluation Program Sand mining' commissioned by the La Mer Foundation (Witbaard et al., 2015). The average water depth was 11 m. 118 119 The bottom frame was located 0.6 km south of the Egmond transect, which was surveyed in 2003, 120 2010 and 2011 to study the location of the turbidity maximum zone (TMZ) along the coast (Van der 121 Hout et al., 2015). The bottom frame was deployed in this TMZ, landward of the average position of 122 the observed peak concentration. One of the measured transects is depicted in Figure 1B to illustrate the spatial context. The study area is located 80 km north of the Rhine outflow in the far-field plume 123 124 of the region of freshwater influence (ROFI) (Simpson et al., 1993). Temporary vertical stratification has been observed at the study site (Rijnsburger et al., 2016; Van der Hout et al., 2015) as well as at a 125 126 location 5 km further offshore (Nauw and Van der Vegt, 2012). The seabed at the study site is sandy 127 with a median grain size of 222  $\mu$ m, the d<sub>50</sub> is 181  $\mu$ m (fine sands), with a clay percentage of 5% in the 128 top layer of 0.05 m (Witbaard et al., 2016; this volume).





Figure 1. A. Map of the research area (Rijkswaterstaat), with the location of the bottom frame (black square) and the two Egmond transects in black lines (dashed: 2003, straight: 2010-2011) near Egmond. The KNMI wind station is in IJmuiden and the Rijkswaterstaat wave buoy IJMSP is indicated by a white square. B. The location of the bottom frame is indicated with reference to a cross-sectional

profile of the SPM concentration along the Egmond transect (August 31, 2011), based on a sequence
of vertical profiles at the vertical dashed lines (Van der Hout et al., 2015).

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### 137 2.1.2 Wind and wave climate

The wind climate at the Dutch coast is characterized by a dominance of southwesterly winds, plus a smaller peak for northerly winds (Figure 2). Both observational years 2011 and 2012 show a stronger dominance of southwesterly winds than in the 10-year average. The wind climate produces a seasonal pattern in the wave climate with a higher wave height and variance in autumn and winter, and lower wave height and variance in spring and summer (Figure 3A).

143 The wave data used in this study are from the wave buoy Munitiestortplaats IJmuiden (IJMSP; 144 Figure 1A). Waves are categorized into storm conditions ( $H_s > 3.5$  m, 5% of time), intermediate conditions ( $H_s = 2 - 3.5$  m, 27% of time) and calm conditions ( $H_s < 2$  m, 68% of time). Storm waves are 145 146 mostly formed by the northwesterly winds (Figure 3B) due to the largest fetch length in the North 147 Sea. Calm waves have a median westerly direction (thus coming equally from the north and south), 148 except for spring when the median direction is NW. The intermediate category waves approach the 149 coast mostly from southerly direction. The wave buoy (35 km offshore) also measured waves from an 150 easterly direction. These wave records are neglected in the analysis, as at 1.5 km from the shore (our 151 study area) such waves have barely developed.

152 Compared to the 10-year average, spring 2011 had milder wave conditions (Figure 3B). In 153 September 2011 ex-hurricane Katia caused wind speeds > 9 Bft and  $H_s$  > 4 m followed by an 154 extremely calm November, and an extremely turbulent December with an average wave height in 155 the intermediate category. Winter 2011-2012 had twice as many storms as normal with four cyclone 156 passings over the North Sea (Roberts et al., 2014), followed by milder wave conditions than normal in 157 February and March 2012.

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Figure 2. Wind roses for the 10-year average (left) and for the years 2011 and 2012. The radial scale of the frequency is 1.4%. Data are provided by KNMI, for location IJmuiden and based on the hourly average.



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Figure 3. Seasonal variation of the wave climate (Data from Rijkswaterstaat). A. The 10-year (2004-2013) monthly average and variance of the significant wave height ( $H_s$ ; calm and intermediate (int.) categories indicated) and the monthly average for 2011 and 2012 (upper panel). B. The monthly median wave direction per wave height category; calm ( $H_s < 2$  m), intermediate (int.;  $H_s = 2 - 3.5$  m) and storm ( $H_s > 3.5$  m) of 10-years of data.

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## 170 2.1.3. Instrumentation

The bottom frame consisted of a triangular aluminum construction with sides of 2 m long and 2 171 172 m high (Witbaard et al., 2015). The frame was equipped with a series of sensors to measure current 173 velocity, turbulence, water depth, temperature, salinity, turbidity and fluorescence. All data was 174 gathered at an interval of 10 minutes. A Nortek Vektor ADV was mounted on the frame at a height of 175 0.3 mab, with the measurement cell at 0.15 mab. The instrument was set to record a burst interval of 176 2 minutes at 1 Hz every 10 minutes. Temperature and salinity were measured with a pumped version 177 of the Sea-Bird SBE 37-SM MicroCAT at 1.4 mab. Turbidity and fluorescence were measured optically 178 with ALEC-JFE Advantech Compact-CLW's at four heights above the bottom, i.e. at 0.3, 0.8, 1.4 and 179 2.0 mab. The sensor at 0.3 mab was the INFINITY version to measure turbidities > 2 gl<sup>-1</sup>. The average 180 of a burst of 10 measurements at 1 Hz was used in the analysis.

181 The deployment period ran from February 25, 2011 until November 24, 2012. The bottom frame 182 was replaced 17 times with intervals between three to ten weeks with a duplicate frame equipped 183 with a clean set of instruments. The duration of the individual deployment periods depended on the season. In spring and summer short periods of three weeks were necessary to limit the effects of biofouling. In autumn and winter, storms sometimes led to longer deployment periods due to operational limits. Besides the short data gaps - up to one day - related to the frame replacements, also longer gaps exist in the data of individual instruments due to loss of battery power when deployment periods became too long or because of malfunctioning of instruments, tilting of the frame or biofouling.

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# 191 **2.2. Data analysis**

The horizontal components of the velocity are rotated with 6 degrees clockwise to align the currents with the coastline, so that  $v_1$  is the alongshore current velocity and  $v_c$  the cross-shore velocity. For some analyses an average was calculated over the tidal cycle, i.e. over the period between two successive slacks after flood.

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### 197 2.2.1. Optical sensor calibration

198 The optical sensors have been calibrated at the start of the project with settled fine sediments 199 which were collected during a previous mooring at the study site. Calibration concentrations ranged from 6 mgl<sup>-1</sup> to 2 gl<sup>-1</sup> for SPM and 0.7 – 475  $\mu$ gl<sup>-1</sup> for CHL-a. The calibration gave a linear response 200 between 20 and 2000 mg l<sup>-1</sup> for SPM ( $R^2 > 0.95$ ), and between 2 and 475 µg l<sup>-1</sup> for CHL-a ( $R^2 > 0.95$ ). 201 202 The used type of optical sensors is known for their stability over long time spans. However, a single 203 calibration cannot account for variations in sediment composition throughout the year or during 204 large resuspension events or varying phytoplankton community compositions. A different type of 205 optical sensor applied in the study site showed a deviation in the calibration of a factor two (Van der 206 Hout et al., 2015).

207

## 208 2.2.2. Grain-size distribution of SPM

At approximately 1 m above the seabed water samples were collected with Niskin bottles during the bottom frame replacement cruises. A 1 l water sample was stored at 4 °C for two days to let sediments settle, then decanted and analyzed in a Beckman Coulter Laser LS 13 320. The samples were only treated with 20 s ultrasonic sound before analysis. Ultrasonic sound may produce air bubbles in the range 100 – 600  $\mu$ m (pers comm. J.B. Stuut), but these were not found in the spectra. Similarly, also a 10 l water sample was stored at 4 °C, decanted and treated in three steps to remove three types of organic components, following the procedure described in Stuut et al. (2014). The first 216 treatment concerned the removal of organic tissue with  $H_2O_2$ . The second and third treatments with 217 HCl and NaOH aimed to remove biogenic calcium-carbonate and biogenic silicates respectively. The 218 remaining fraction is called the insoluble or terrigenous fraction. Visual inspection of microscope 219 images was done to ensure complete removal of the organic particles. The samples were split into 220 two subsamples before the treatments, one subsample for the analysis of the grain-size distribution 221 and one subsample to determine the mass changes per treatment. For the mass analysis the 222 subsamples were centrifuged and freeze-dried after treatment and then weighed to determine the 223 reduction in weight.

224

#### 225 2.2.3. Bed shear stress from ADV

For the calculation of the bed shear we followed the procedure in Verney et al. (2007) with the exception that wave shear stress was determined according to the method of Van Rijn (2007). The latter gives an all inclusive bed roughness predictor for a combination of currents and waves with all bed forms in the field. The mean bed shear stress is calculated with (Verney et al., 2007);

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$$\tau_{\rm m} = \tau_{\rm c} \left( 1 + 1.2 \left( \frac{\tau_{\rm w}}{\tau_{\rm w} + \tau_{\rm c}} \right)^{3.2} \right) = \tau_{\rm TKE} = 0.5 \rho C \sqrt{u_t'^2 + v_t'^2 + w_t'^2}$$
(1)

with  $(u_t^{\prime 2},v_t^{\prime 2},w_t^{\prime 2})$  calculated from the spectral analysis of the fluctuating part of the velocity 231 232 (u', v', w') (Soulsby and Humphery, 1990). The constant C has been found to approximate 0.19 for a 233 wide variety of flows (Verney et al., 2007). The fluctuating velocities are first despiked with the 234 phase-space thresholding method developed by Mori et al. (2007) and Goring and Nikora (2002). The 235 method requires a longer time series than the burst length to separate the spikes from the standard deviation. Therefore all data is merged and the despike method is applied on subsets with a length of 236 237 250.000 data points. Subsequently the data are separated back into burst intervals. Spectra are 238 obtained from the despiked burst intervals with the fast-fourier transform, and are averaged over an 239 hour. In the logarithmic space of the (averaged) spectra, the separation between the turbulent 240 velocities and the wave orbital velocities is drawn as a linear line between the (non-logarithmic) 241 frequencies 0.04 and 0.31 Hz. This range of wave frequencies covers both swell and wind waves 242 (bimodal wave spectrum) and is well outside the influence of the Nyquist frequency ( $f_N = 0.5$  Hz). The two resulting surface areas of the linear spectra are  $u_t'^2$ ,  $v_t'^2$ ,  $w_t'^2$  (turbulent variance) and  $u_w'^2$ ,  $v_w'^2$ ,  $w_w'^2$ 243 244 (wave variance). When no waves are present, the separation line follows the inertial slope of the 245 spectrum. The 1 Hz frequency and the 120 s burst length of the measurements proved to be just high 246 and long enough - when averaged over one hour - to distinguish the ranges of the wave orbital 247 velocities from the inertial subrange, and thus to separate the turbulent and wave signals from each 248 other.

(4)

249 An important parameter in the calculation of the wave shear stress is the bed roughness ( $k_s$ 250 or  $z_0$ ), as its value can vary within 2 orders of magnitude. The bed roughness depends on the bed 251 form, which may alter quickly when conditions change, and is thus a sensitive parameter. After evaluating various bed roughness predictors for a similar coastal site off Terschelling (The 252 Netherlands, water depth 5 - 10 m, sand of 200 µm, orbital velocities up to 0.6 ms<sup>-1</sup>), Houwman 253 254 (2000) advised a constant bed roughness ( $k_s$ ) of 0.1 m. He stated that a uniform bed roughness was 255 the optimum estimate for all flow conditions, with a similar accuracy as the uniform bed roughness 256 predictor by Van Rijn (2007).

257 The wave friction factor  $f_w$  is then calculated as (Van Rijn, 2007)

258

$$f_{w} = \exp\left(-6 + 5.2 \left(\frac{U_{w} T_{p}}{2\pi k_{s}}\right)^{-0.19}\right)$$
<sup>(2)</sup>

With  $U_w$  and  $T_p$  determined from the ADV. The wave period  $T_p$  was determined from the frequency of the highest spectral peak from the pressure sensor. The representative bottom orbital velocity ( $U_w$ ) is determined from the wave variance of the velocity spectrum, with (Wiberg and Sherwood (2008))

263 
$$U_{w} = \sqrt{2} \cdot \sqrt{u_{w}^{\prime 2} + v_{w}^{\prime 2}}$$
(3)

The wave shear stress ( $\tau_w$ ) and the wave-current shear stress ( $\tau_{wci}$ ) are then obtained with (Verney et al., 2007):

$$\tau_w = 0.5 \rho f_w U_w^2$$

267

7 
$$\tau_{wci} = \tau_{max} = \sqrt{(\tau_m + \tau_w \cos(\Psi))^2 + (\tau_w \sin(\Psi))^2}$$
 (5)

268 With  $\Psi$  the angle between the wave direction (from wave buoy) and the current direction ( $\overline{u}, \overline{v}$ ). 269

## 270 2.2.4. Rouse number

The Rouse profile describes in a simple way the logarithmic distribution of the SPM over depth by capturing the vertical profile in one number; the Rouse number  $R_n$ . The Rouse equation is based upon the equilibrium state of the concentration, in which the vertical upward flux by turbulence and the gravitational settling are in balance (Rouse, 1950). It is assumed that gradients in the advective (i.e. horizontal) transport are much smaller than in the vertical and therefore negligible. The common representation of the Rouse equation is (Mehta, 2014)

277 
$$\frac{C(z)}{C_a} = \left[\frac{z_a(h-z)}{z(h-z_a)}\right]^{W_s/\kappa u_*}$$
(6)

278 With the reference concentration  $C_a = C(a)$  at reference height  $z = z(a) = z_a$ , h = water depth,  $w_s =$ 279 settling velocity,  $\kappa$  = Von Karman constant and  $u_x$  = the friction velocity. The eddy diffusivity is 280 described with a parabolic form, which assumes the absence of stratification. The settling velocity is assumed to be independent of the sediment concentration and the shear rate. The dimensionless ratio in the exponent of equation (6) is the Rouse number  $R_n = w_s/\kappa u_*$  (Rouse, 1950). It represents the ratio between the downward motion of the sediment by gravity and the upward force by the turbulent motion. The value of  $R_n$  gives an indication of the suspension mode; a high  $R_n$  represents a steep concentration gradient, and a low  $R_n$  a nearly uniform profile. The separation between both modes is usually drawn at  $R_n = 0.8$  (Mehta, 2014).

The Rouse number is calculated from an observed profile by rewriting equation (6) (Mehta, 2014):

$$R_n = \frac{-\log(C/C_a)}{-\log((z_a (h-z))/z(h-z_a))}$$
(7)

290 The reference concentration  $C_a$  is taken at  $z_a = 0.3$  mab: the observation closest to the seabed. For the calculation of  $R_n$ , the SPM and CHL-a concentrations have been smoothed with a low pass 291 292 band filter with a 50 minutes window. The nominator and the denominator on the right-hand side of 293 equation (7) were calculated for each of the 3 remaining heights z = [0.8, 1.4, 2.0] mab and with h =294 11 m.  $R_n$  is the coefficient of the linear regression without offset through these three points if  $R^2 > 1$ 295 0.95. If R<sup>2</sup> is less than 0.95, we assume that the vertical concentration distribution does not follow a 296 Rouse profile.  $R_n$  could not be calculated for the period between October 2011 and early January 2012 since the optical sensor at 2.0 mab malfunctioned in this period, and not enough information 297 298 was available for a reliable regression.

299

# 300 *3. Results*

# 301 *3.1. Hydrography*

#### 302 3.1.1. Local environmental conditions

303 The study area shows a seasonal cycle in the water temperature (Figures 4F and 5F) ranging from 304 3 °C in February 2012 to 20 °C in August 2012. A tidal fluctuation of maximum 1 °C is sometimes 305 apparent during calm conditions such as in November 2011 and February 2012. The salinity is 306 relatively low due to the large Rhine river outflow and varies between 25 and 32 PSU, generally with values close to 30 PSU (Figures 4G and 5G). Fluctuations with a magnitude of maximum 2 PSU occur 307 308 on a tidal scale. There is no distinct seasonal pattern in salinity but larger fluctuations occur on time 309 scales of weeks to months. Lowest salinities occurred in January 2012. The tidal range varied over the 310 spring-neap cycle between 1.1 and 1.9 m (Figures 4E and 5E). The near-bed alongshore velocity amplitude ranged between 0.5 and 0.2 ms<sup>-1</sup> over the spring-neap cycle. The spring-neap cycle is 311 312 clearly present in the water level signal, but is distorted in the velocity signal by intermediate and 313 storm waves.



- Figure 4 (A-I). Data recorded at the bottom frame in 2011. SPM concentrations at the lowest 315 316 observation (0.3 mab, panel A) and at 1.4 mab (panel B) including the tide-averaged value (black 317 lines). CHL-a concentrations at the lowest observation (0.3 mab, panel C) and at 1.4 mab (panel D) 318 and their tide-averaged value (black lines). Panel E shows the tidal amplitude of the water level and 319 the alongshore tidal current. The moon phases (full moon;  $\circ$ , and new moon;  $\bullet$ ) are shown; spring tide is two days after full and new moon. Panel F shows the temperature signal, and the deployment 320 321 and recovery dates as well. Panel G shows the salinity. Both the bed shear stress ( $\tau_{MAX}$ ; blue) and the 322 significant wave height (H<sub>s</sub>; black, source; Rijkswaterstaat) are shown in panel H. The Rouse number
- 323  $R_{n}$ , and the percentage of fitted Rouse profiles per tidal cycle,  $Q_{p}$ , are shown in panel I.



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#### 327 3.1.2. Grain size suspended particles

328 Only five water samples contained enough SPM to measure the grain size of the terrigenous 329 sediment fraction. The grain-size distributions (Figure 6; solid lines) of the terrigenous particles entail 330 the clay fraction (< 2  $\mu$ m) and the silt fraction (< 63  $\mu$ m). Peaks in the distributions are at 0.4  $\mu$ m, 5 331  $\mu$ m, 10  $\mu$ m and 30 - 50  $\mu$ m, which are present regardless of the time in the year. The mode of the 332 terrigenous particles is at  $10 \pm 5 \mu m$  (Figure 6, solid lines). The modes of the untreated (but sonified) samples have the same values (dashed lines in Figure 6). However, the tails of the bell-shaped 333 334 distribution are different before and after treatment. All samples except of Feb 2012 have particles in 335 the range 50 - 200  $\mu$ m, which is the typical range for phytoplankton; both biogenic silicates and 336 foraminifera (Stuut et al., 2002). Notably, the clay fraction is more present in the treated samples 337 than in the untreated. An explanation could be that before the treatment this fraction was still 338 captured in flocculi with a size similar to silt, despite the high turbulence in the measurement volume 339 of the Coulter Laser (the exact value is not calculated). After chemical treatment these flocculi have 340 disintegrated into their primary small particles, pointing at a strong aggregation of the clay particles.

The mass analysis shows that the SPM consists for 35 - 55% of organic particles, being: organic tissue (10 – 15%), biogenic silicate (0 – 20 %) and biogenic calcium-carbonates (CaCO<sub>3</sub>; 25 – 30 %). A seasonal trend analysis is not possible due to the limited amount of samples (5). The high organic content in the samples suggests that at our study site the terrigenous particles are mostly present in aggregated form, as observed by others in similar environments (Fettweis et al., 2014; Jones et al., 1998).



Figure 6. The grain size distribution of five samples before treatment (dashed lines) and after
 treatment (solid lines)

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#### 352 *3.1.3. Bed shear stress*

The bed shear stresses varied between 0.1 and 40 Nm<sup>-2</sup> (Figures 4H and 5H) and has a wave and a tidal component. The bed shear stress exerted by waves is 1 - 2 orders of magnitude larger than the range due to the tidal variation which is between 0.2 and 1 Nm<sup>-2</sup> during calm conditions (Figures 4H and 5H). The scatter plot in Figure 7 indicates a positive logarithmic relation between the bed shear stress and the significant wave height for  $H_s > 0.8$  m. The variation around the median is due to the tidal variation.



Figure 7. Scatter plot of the bed shear stress ( $\tau_{MAX}$ ) versus the significant wave height ( $H_{s}$ ; 1 hr intervals). The large white dots are the median bed shear stresses per 0.2 m  $H_s$  bins.

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# 363 3.2. SPM and CHL-a variation

## 364 3.2.1. SPM concentrations

365 Panels A and B in Figures 4 (for 2011) and 5 (for 2012) show the SPM concentrations at 0.3 mab 366 and at 1.4 mab. The SPM concentrations ranged from lower than 20 mgl<sup>-1</sup> and attained at least 2 gl<sup>-1</sup> 367 (the limits of the linear gain of the optical sensors). The time series shows fluctuations of the SPM 368 concentration on both the short term (tidal scale, regular) and the long term (storms, irregular). The 369 tidal SPM variation (panels A and B) increases during intermediate ( $2 < H_s < 3.5$  m) and storm ( $H_s > 3.5$ 370 m) conditions and decreases afterwards during calm conditions (H<sub>s</sub> in panels H). The peaks in the 371 SPM concentration occur at an irregular interval and seem related to elevated bed shear stresses due to increased wave heights (both Figures 4H and 5H). At 0.3 mab the SPM concentrations in these 372 373 peaks range from between 0.5 gl<sup>-1</sup> to at least 2 gl<sup>-1</sup>. In spring and summer, peaks barely exceed 1 gl<sup>-1</sup>, while in the 2011-2012 winter all five storms caused elevated SPM concentrations exceeding 2 gl<sup>-1</sup>. 374 Even the SPM concentrations at 1.4 mab reached up to 1  $gl^{-1}$  in December 2011, far higher than the 375 average peaks of 0.5 gl<sup>-1</sup> during the remainder of the year. January and February 2012 remained 376 377 turbid with concentrations regularly exceeding 1 gl<sup>-1</sup> at 0.3 mab, while spring was a calm period with 378 relatively low peak concentrations of around 0.5 gl<sup>-1</sup>. In the summer of 2012 concentrations regularly 379 exceeded 1 gl<sup>-1</sup>, and in autumn two storms ( $H_s > 3.5$  m) passed during our deployments, but 380 unfortunately the SPM concentrations were not recorded due to failing devices. Later in autumn 381 concentrations exceeded 2  $gl^{-1}$  a few times at 0.3 mab.

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#### 383 3.2.2. CHL-a concentrations

384 Panels C and D in Figures 4 (for 2011) and 5 (for 2012) show CHL-a concentrations between 3  $\mu$ gl<sup>-1</sup> to at least 300  $\mu$ gl<sup>-1</sup> at 0.3 mab and 1.4 mab. The concentration at 0.3 mab is always higher than 385 386 at 1.4 mab. Both the relatively high concentrations (the order of surface concentrations is 10  $\mu$ gl<sup>-1</sup> 387 (Joordens et al., 2001)) and the increase towards the seabed illustrate that the observed CHL-a at 0.3 388 mab is an accumulated suspension near the seabed. The CHL-a concentration shows variations on 389 both the regular tidal and the irregular storm scale, similar to what was observed for the SPM 390 concentrations and most likely related to resuspension of the accumulated suspension. The 391 dominant variation in the CHL-a concentration is the increase in spring and decrease in summer: the 392 spring bloom. There are differences in the development and magnitude of the spring bloom of both 393 years. In 2011 the CHL-a concentration at 0.3 mab built up in April during calm conditions, and the 394 near-bed peak occurred half of May with tide-average concentrations of  $125 \,\mu gl^{-1}$  and peaks of > 300 µgl<sup>-1</sup>. In 2012 the CHL-a concentration started to build up at the end of March, but stopped at the 395 end of April. Concentrations had reached near-bed peaks of 200 µgl<sup>-1</sup> and tide averaged 396 397 concentrations 100 µgl<sup>-1</sup>, somewhat lower than in 2011. In April 2012 several periods characterized 398 by intermediate wave conditions occurred which seem to have suppressed the further development 399 of the phytoplankton bloom. Figure 3A suggests that April 2012 was a period with higher wave 400 energy input and deviates from the long term average conditions. Half of May 2012 concentrations 401 increased again to similar tide-averaged concentrations as in April. Other, short-lived peaks in the 402 CHL-a concentration at 0.3 mab occurred at moments of increased bed shear stresses at the same 403 instant that SPM concentrations increased. These peaks can be up to 100 µgl<sup>-1</sup>, even in winter. Tideaveraged concentrations at 1.4 mab attained 100  $\mu$ gl<sup>-1</sup> in spring 2011 and 40  $\mu$ gl<sup>-1</sup> in spring 2012, 404 405 which are only slightly lower than at 0.3 mab. Conversely, the incidental wave-related peaks at 1.4 406 mab are an order of magnitude lower and rarely exceed 15  $\mu$ gl<sup>-1</sup>.

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### 408 *3.2.3. The seasonal scale*

409 Figure 8 shows the GAM smoothed yearly variations of SPM and CHL-a at 0.3 mab on the basis of 410 weekly average concentrations (Generalized Additive Model; Wood (2006)). The GAM model estimates a smoother function describing the seasonal variation for each year. The smoother 411 412 functions for each variable (CHL-a, SPM) and each year were significant. This shows that there is a 413 seasonal variation in the near-bed concentrations of both CHL-a and SPM. At the same time the 414 smoother functions for each variable are compared for a difference between years. This comparison 415 showed that the seasonal development of near-bed CHL-a concentrations differed between 2011 and 416 2012. Such a seasonal difference could not be demonstrated for the near-bed SPM concentrations. 417 The near-bed CHL-a concentrations show a sharp peak in 2011 and a wide peak in 2012 (Figure 8A). 418 In the summer period of 2011 the near-bed SPM concentrations show minimum values but in 2012 419 these minimum values are overshadowed in June (Week 26; Figure 8B) when intermediate waves hit 420 the Dutch coast (Figure 5H).

Winter weekly-averaged SPM concentrations reach 250 mgl<sup>-1</sup>, summer values are between 50 and 150 mgl<sup>-1</sup>. The near-bed CHL-a increase related to the spring bloom starts in week 11 (2011) or week 6 (2012) when average concentrations increase from lower than 10  $\mu$ gl<sup>-1</sup> to 30 - 60  $\mu$ gl<sup>-1</sup> and decrease at the end of May (week 20). The autumn and winter values remain elevated around 10  $\mu$ gl<sup>-1</sup> and show oscillations with the wave climate. The spring bloom related CLH-a increase in 2012 seems to start in February (Figure 8A), but from the raw time-series in Figure 6C it follows that the
CHL-a increased temporarily in February and March due to two intermediate wave events, and the
steep increase started at the end of March.

429





Figure 8. The GAM modeled seasonal cycles in the CHL-a (panel A) and SPM (panel B) concentrations at 0.3 mab over the period February 2011 – November 2012 on basis of weekly averaged concentrations. The shaded areas provide the 95% confidence around the estimates (Wood 2006). The seasons are indicated at the top of the graphs. In September 2012 (week 35 – 39) no (concentration) data was recorded, but estimated by the model. Therefore the confidence range is larger in this period.



#### 438 *3.2.4. The tidal scale*

In Figure 9 the power spectral density (PSD) estimates for the SPM and CHL-a concentrations at 439 440 0.3 and 1.4 mab are shown, besides those for the current velocity and the water level. The current 441 velocities and the water level show peaks at the lunar tidal frequencies. The focus on the M<sub>2</sub> and M<sub>4</sub> 442 frequencies in Figure 9B reveals a contribution of the  $S_2$  tide, though an order of magnitude smaller. 443 The interaction of the M<sub>2</sub> and S<sub>2</sub> frequencies generates the spring-neap cycle. The PSDs of the SPM 444 and CHL-a concentrations show also peaks corresponding to the lunar tidal frequencies. The PSD 445 estimate of the SPM concentration at 0.3 mab has three peaks at frequencies corresponding to the 446 tidal frequencies M<sub>2</sub>, M<sub>4</sub> and M<sub>8</sub>. At 1.4 mab similar peak frequencies, but lower and broader spectral 447 peaks occur. In SPM and CHL-a, M<sub>4</sub> is particularly strong, an indication of the semi-diurnal resuspension (i.e. with each ebb and flood). The PSD of the CHL-a concentration at 0.3 mab and at 448 449 1.4 mab has four peaks at frequencies corresponding to the tidal frequencies M<sub>2</sub>, M<sub>4</sub>, M<sub>6</sub> and M<sub>8</sub>. 450 Compared to the SPM concentration, the CHL-a concentration oscillates at an additional frequency 451 i.e. the M<sub>6</sub> tidal frequency. The largest spectral peak of the CHL-a concentration at 0.3 mab is at the 452  $M_4$  frequency, while at 1.4 mab the  $M_2$  equals the  $M_4$  peak. As the spectral peaks of the SPM and 453 CHL-a concentrations are already broader than those of the velocities and water level, it is difficult to 454 distinguish their spring-neap oscillation.

455 Figure 10 gives three examples of the tidal variation in the concentrations of SPM and CHL-a 456 during two calm conditions and during a storm. The left panels are from October 2nd-3rd 2011 after 457 a long period of calm weather. The middle panels are at the end of a storm ( $H_s \sim 3 - 4$  m). The right 458 panels represent October 14th-15th 2011, 2 days after a 6-day storm period. In the left and right-459 hand panels the bed shear stress and the wave height have a similar magnitude, but the tidal 460 oscillation of the SPM concentration after the storm is one order of magnitude larger than before. 461 The range of the intratidal SPM variation appeared quite large with an average of 60 mgl<sup>-1</sup> before the storm and up to 400 mgl<sup>-1</sup> after the storm. Notably, during the storm still a tidal signal is apparent, 462 463 but less clear. During both calm periods 8 SPM concentration peaks appear at the 8 velocity peaks (4 464 flood and 4 ebb) corresponding to the  $M_4$  spectral peak in the concentration and the  $M_2$  spectral 465 peak in the velocity. The most obvious relation of the concentration signal with the tide forms the exact correspondence in timing of the SPM and CHL-a minima and the slack tides (i.e. minimal 466 467 velocities). On the other hand, the timing of the concentration peaks has a phase difference with the 468 tidal current peaks, varying per tide, showing concentration peaks either before or after the 469 maximum current velocity. Sometimes the concentration maxima are double-peaked.

During the calm period the CHL-a concentration oscillation with the tidal current is smaller than the SPM oscillation, with 4 CHL-a peaks barely noticeable for 8 velocity peaks (Figure 10C). During and after the storm period the CHL-a concentrations show exactly the same signal as the SPM concentrations, as if being merged together (Figure 10G). Figure 10 suggests that the contribution of CHL-a to the SPM is different before and after the storm.







Figure 9. Power spectral density (PSD) estimates of the SPM concentrations, the CHL-a 477 478 concentrations, the water level and the current velocity; both alongshore  $(v_l)$  and cross-shore  $(v_c)$ . The PSDs are calculated with the Welch method with a window length of 4000 over the period March 1st 479 480 2011 – July 1st 2012. The main lunar tidal frequencies ( $M_2$ ,  $M_4$ ,  $M_6$  and  $M_8$ ) and  $S_2$  are indicated by 481 vertical dashed lines. Indicated as well is the -5/3 slope (black dashed line), which is indicative for the 482 inertial subrange. On the right is a zoom on the  $M_2$ ,  $S_2$  and  $M_4$  frequencies of all variables.



484

Figure 10. The variation of the SPM and the CHL-a concentrations over 48 h after a calm period 485 486 (panels A-D; October 2-3, 2011), during a storm (panels E-H; October 10-11) and after a storm period (panels I-L; October 14-15, 2011). Panels A, E and I show the SPM concentrations at 0.3 mab (black) 487

and 1.4 mab (grey) during all periods. Panels B, F and J give the alongshore current velocities. Panels
C, G and K show the CHL-a concentrations at 0.3 mab (black) and 1.4 mab (grey). Panels D, H and L
give the variation of the bed shear stress (τ<sub>MAX</sub>; black line; logarithmic scale) and the significant wave
height (H<sub>s</sub>; grey line; linear scale).

492

# 493 *3.2.5. The effect of the wave climate*

In Figure 11 the SPM concentration at 0.3 mab is plotted against the bed shear stress. The scatter 494 495 plot shows a large spread of the data, and the tidal variation of the bed shear is visible for  $\tau_{MAX} = 0.2$  -496 1 Nm<sup>-2</sup>. Three observations can be drawn from the relation between bed shear stress and the near-497 bed SPM concentration. The first is a positive relation between the bed shear stress and the SPM 498 concentration; i.e. increasing bed shear stresses due to waves lead to an increase in the SPM 499 concentration. The second observation is that the median concentration has a peak at a shear stress 500 of 6  $\text{Nm}^{-2}$  (coinciding with median H<sub>s</sub> = 2.5 m). And lastly, peak SPM concentrations (> 1 gl<sup>-1</sup>) do not 501 occur at maximum shear stresses but at shear stresses between 0.4 and 15 Nm<sup>-2</sup>. The explanation for 502 these last two observations is shown in the middle panel of Figure 10. Peak SPM concentrations arise 503 in the aftermath of the storm, at 18:00 on October 11, while during the storm the SPM concentration 504 has lower values. The decreasing shear stress initiates sinking of SPM into a high concentration 505 suspension layer near the seabed. An analysis of the wave direction showed only a marginal effect of 506 the wave direction on the magnitude of the SPM concentration.



507

508 Figure 11. Scatter plot of the SPM concentration and the bed shear stress at 0.3 mab. The large 509 red dots are the median SPM concentrations.

# 510 **3.3.** The vertical profile

#### 511 3.3.1. Frequency distribution

The SPM and CHL-a concentrations have a skewed distribution to the left at all four heights, and 512 513 therefore their frequency distributions are described by both the lognormal and the exponential 514 distribution (Figure 12). The geometric mean and standard deviation (lognormal distribution) and rate parameter (exponential distribution) are given in Table 1. As the dataset covers one year and 515 516 nine months (one winter is missing), the dataset was split in two overlapping year-long time series: 517 March 2011 - March 2012 and July 2011 - July 2012. There was no bias found for the resulting 518 dataset length regarding the SPM concentrations since the average and standard deviation of the 1.8 519 year dataset were similar to those of the two subsets. For both variables the geometric mean 520 decreases from the seabed upward. For the SPM concentrations the decrease is logarithmically from 521 83 mgl<sup>-1</sup> to 45 mgl<sup>-1</sup>. However, the geometric mean of the CHL-a concentrations decreases only 522 between 0.3 and 0.8 mab from 13.7 to 9.2 µgl<sup>-1</sup>, and remains constant between 0.8 and 2.0 mab.



525 Figure 12. Frequency distributions of the SPM (A) and CHL-a (B) concentrations at 0.3 m, 0.8 m, 526 1.4 m, and 2.0 m above the seabed, both the histogram and the fitted exponential function.

527

529 ( $\lambda$ ) of the exponential distribution of the SPM (in mgl<sup>-1</sup>) and the CHL-a (in  $\mu$ gl<sup>-1</sup>) concentrations at all

530 *four heights.* 

	SPM			CHL-a		
	geo. mean	geo. std	λ (mg l <sup>-1</sup> )	geo. mean	geo. std	λ (μg l <sup>-1</sup> )
	(mg l <sup>-1</sup> )			(µg l <sup>−1</sup> )		
2.0 mab	45	2.4	59	9.1	2.3	11
1.4 mab	49	2.6	66	8.7	2.2	11
0.8 mab	60	2.7	80	9.2	2.3	12
0.3 mab	83	3.3	106	13.7	2.4	18

531

# 532 *3.3.2 Rouse numbers and profiles*

For 84% of the profiles with SPM concentrations between 20 and 2000 mgl<sup>-1</sup> the fit had an R<sup>2</sup> 533 534 larger than 0.95 and thus followed a Rouse profile from which a reliable  $R_n$  could be calculated. 535 About 26% of the profiles had a concentration either below 20 mgl<sup>-1</sup> or above 2000 mgl<sup>-1</sup> and were 536 omitted from the analyses. These concentrations could not reliably be detected by the sensors and 537 thus cannot give a reliable estimate of  $R_n$ . The  $R_n$  was positively correlated with the concentration at 0.3 mab ( $R^2 = 0.6$ ). The CHL-a concentration profiles fitted the Rouse profile only for 2% of the data, 538 539 and are therefore not included in further analyses. This result confirms the result in section 3.3.1 540 concerning the difference between the uniform CHL-a distribution upward of 0.8 mab and the 541 logarithmic SPM distribution.

542 The calculated Rouse number for the SPM profiles varied over time both at the tidal and the episodic storm scale between 0.01 and 2.5 (Figures 4I and 5I) which suggests a wide variety of SPM 543 544 concentration profiles. By extrapolating the near-bed measurements up to the surface following the Rouse profile, the SPM profiles can be divided into a near-bed suspension ( $R_p > 0.8$ ) and a mixed 545 546 suspension ( $R_n < 0.8$ ) (Mehta, 2014). In 90 % of the time  $R_n$  is found to be smaller than 0.8, and thus 547 the SPM is dominantly present in a mixed suspension, ie. the SPM profile is more uniform up to the 548 water surface. In a small percentage of the time (10%) the SPM suspension is concentrated in the 549 near-bed layer (near-bed suspension).

The temporal variation in the percentage of SPM profiles following a rouse profile ( $R^2 > 0.95$ ) is also shown in Figures 4I and 5I and quantified as the parameter  $Q_p$ , which is the percentage of wellfitted Rouse profiles per tidal cycle. The average value of  $Q_p$  is 88%, with low  $Q_p$  predominantly apparent in March 2011, and February / March 2012. It appeared that SPM profiles with a reference concentration (at  $z_q = 0.3$  mab) larger than 500 mgl<sup>-1</sup> always followed a Rouse profile. To explore whether the fit of the Rouse profile to the SPM profile depended on the tidal phase, the percentage of accepted and rejected Rouse profiles is plotted against the alongshore current velocity (Figure 13). This figure is based on the tides with less than 20% rejected Rouse profiles ( $Q_{\rho} >$ 80%) to eliminate causes of potential deviations other than the tidal flow. From the figure it is evident that during all phases of the tidal cycle Rouse profiles are rejected (blue bars in Figure 13B). The highest percentage of rejected Rouse profiles is found at low current velocities and more precisely at the accelerating flow after slack tide up to 0.5 of the maximum tidal velocity (Figure 13A).



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Figure 13. A: Percentage of rejected Rouse profiles per phase of the tide. The latter is calculated as the alongshore velocity divided by the maximum alongshore velocity of the tidal cycle. The graph shows the percentage rejected of all tides with  $Q_p > 80\%$ , and the percentages when the velocities are separated in accelerating (acc) and decelerating (dec) tidal currents. B: The raw numbers of the rejected observations and all selected observations per tidal phase.

## 568 **4. Discussion**

## 569 **4.1. Seasonality**

570 The long-term signals of the near-bed SPM and CHL-a concentrations show opposite seasonal 571 cycles in the Dutch coastal zone (Figure 8). A seasonal cycle in the SPM concentrations (high in 572 winter, low in summer) has also been reported for similar environments, such as the Irish Sea (Jago et al., 2007) and other parts of the southern North Sea (Blauw et al., 2012; Fettweis et al., 2012). In 573 574 these studies the wave climate was the main driver of the seasonal SPM variation (Stanev et al., 575 2009). Our study confirms this larger role of the wave climate, since the bed shear stress - the 576 principal parameter for SPM resuspension – is logarithmically related to the wave height and reaches 577 values up to two orders of magnitude larger for wave heights than the tidal current (Figure 7).

578 Recently however, Fettweis et al. (2014) suggested that the seasonal variation in SPM is 579 primarily governed by the seasonal variation in organic matter, and suggested the wave climate as a 580 secondary process. They presented a seasonal variation in the floc aggregation due to the variation in 581 organic matter, hence in fall velocity, leading to a seasonal signal in the settlement of SPM. It is 582 indeed conceivable that after the initial decrease of SPM concentrations, the first phytoplankton 583 bloom initiates a possible feedback loop, in which the freshly formed organic matter facilitates SPM 584 settlement, thereby improving light conditions even more and so on, as suggested by Blauw et al. 585 (2012). Our summer SPM concentrations give a hint of the interference of organic matter, as in both 586 years the summer SPM concentrations were lower than in winter for similar wave heights. Not only scavenging by sinking aggregates (Blauw et al., 2012; Jago et al., 2007; Fettweis et al., 2014), but also 587 588 stabilization of fines in the seabed by the dominant Ensis Directus may have reduced the 589 resuspension signal in summer (Borsje et al., 2008; Witbaard et al., 2016; this volume).

590 The CHL-a signal in our study shows an increase in spring - when SPM decreases-, but our 591 results cannot point to a leading role for CHL-a in the seasonal pattern of SPM. A translation from 592 CHL-a concentrations to phytoplankton species, to excretion of (sticky) polysaccharides, to aggregate floc sizes and finally SPM settlement cannot be made. Also, we do not know if and how the spring 593 594 increase in the near-bed accumulated CHL-a signal is representative of the phytoplankton growth 595 higher up in the water column, or to what extent a phase lag exists at this shallow site. Our results do 596 point to a principal role for the wave climate in the seasonal variation of SPM. The oscillations in the 597 seasonal SPM concentration show a good agreement with the different variations in the wave 598 climate in both years.

## 599 **4.2 Fluffy layer development**

Waves increased the tidal bed shear stress range of 0.2 - 1 Nm<sup>-2</sup> with one to two orders of 600 magnitude up to 40 Nm<sup>-2</sup> (Figure 7). These high bed shear stresses mobilize the seabed and 601 602 resuspend large amounts of sediment. Remarkably, the maximum SPM concentrations exceeding 1 603  $gl^{-1}$  were not measured at the peak of the storms (Figure 11), but arose in the aftermath (Figure 10E). 604 The relative low concentrations during the heaviest storms may have two reasons. Firstly, they may 605 have been caused by the low sensor sensitivity for sand grains. Optical backscatter is known to be 606 sensitive for the finest particles and thus not fully representative for the resuspension of the grain 607 size range of the seabed (Connor and De Visser, 1992). Another likely explanation is however the 608 unbalance between the release rate of fines during the development of a storm and the high sinking 609 rates in the aftermath of a storm.

610 Over the course of a storm development (buried) fines are gradually resuspended from 611 deeper sediment layers and are thus slowly released. At increasing bed shear stresses fines from 612 deeper and deeper layers are "washed out" and mixed over the entire water column where it forms 613 a diluted suspension. Once the storm has passed and wave action decreases, the sediment load from 614 the entire water column starts to settle and accumulates in a confined layer near the seabed: an 615 increase in SPM concentration at 0.3 mab at low turbulence levels (Figure 10E). The high Rouse 616 numbers for high concentrations (linear relation  $R_a$  with  $C_a$ ) indicate a steep vertical concentration 617 gradient, which supports the development of a highly concentrated fluffy layer close to the seabed 618 after a storm.

This (fluffy) layer also contains high concentrations of CHL-a rich particles that settle to the seabed over time. The settlement, subsequent burial and following resuspension of CHL-a concentrations is not uncommon and has been observed before (Boon and Duineveld, 1998; Jago et al., 2007). Such settlement leads to higher levels of near-bed CHL-a concentrations than surface water concentrations (order of 10  $\mu$ gl<sup>-1</sup> (Joordens et al., 2001)). The high abundance of CHL-a in the study area close to the seabed is a possible explanation for the high abundance of the benthic filter feeder *Ensis directus* in the area (Witbaard et al., 2015).

## 626 4.3 Tidal scale resuspension

The spectral analysis (Figure 9) and the 48 hour zooms (Figure 10) indicated that SPM and CHL-a concentrations have a tidal oscillation. Single point observations always make it difficult to assess whether SPM variations are caused by local resuspension and settling processes or by horizontal advection, as both processes create similar concentration signals (Bass et al., 2002; Blewett and Huntley, 1999; Jago et al., 2006). In our study area both processes may act, as the requirements for both are present: 1) a horizontal gradient in the SPM concentration (Figure 1) and 2) a local source (fluffy layer above the sandy seabed).

634 The spectral analysis (Figure 9) identified M<sub>4</sub> as the main tidal component in both the SPM and 635 the CHL-a dataset, followed by the M<sub>2</sub> component. The spring-neap cycle could not be distinguished 636 in the signal due to the flattening of the spectral peaks: a result of the phase differences of the peak 637 concentrations with respect to the velocity peaks. The presence of the M<sub>4</sub> signal can be due to both 638 the resuspension and settling process and the advection of a concentration gradient back and forth 639 along the bottom frame with the M<sub>2</sub> tide (Blewett and Huntley, 1998). The advective component may 640 be detected as a phase difference between the peaks of the concentration and the velocity, whereas 641 the resuspension process has no phase differences - when measured close the seabed (Bass et al., 642 2002). As the troughs in the SPM concentration time-series coincided clearly with the slack periods of 643 the tidal cycle, we suggest resuspension takes up a larger role than advection.

The M<sub>2</sub> component in the concentration signals may have multiple causes, amongst others a difference in resuspension strength by the ebb (weaker) and flood (stronger) currents, but also the sensitivity of the optical sensor to the aggregation and break-up process on the tidal cycle has been observed (Jago et al., 2006). The latter produces a typical 'twin-peak signature' in the concentration time-series of which a modified version can be observed in Figure 10 before and after the storm.

#### 649 **4.4 Cross-shore advection**

Advective transport can occur both in the alongshore and the cross-shore direction. In both 650 651 directions the velocity spectra displayed  $M_2$  as the principal tidal component (Figure 9). In the 652 alongshore direction the velocities are largest (Figures 4E and 5E) but the concentration gradient is 653 small and not accurately known (Van der Hout et al., 2015). In the cross-shore direction the velocities 654 are small but the concentration gradient is large (Figure 1). Two processes can generate a cross-shore 655 oscillation on the tidal scale; the tidal ellipse (a continuous process; Houwman, 2000) and tidal 656 straining (an infrequent process; Simpson et al., 1993). As the discrimination between both processes 657 is a study in itself, we will here only focus on the measured quantities.

The question is how much the cross-shore advection may contribute to the tidal SPM concentration variation measured at the bottom frame: on average 84 mgl<sup>-1</sup> at 1.4 mab and 174

660 mgl<sup>-1</sup> at 0.3 mab. By combining SPM concentration data obtained from the spatial observations in 661 2010-2011 (Van der Hout et al., 2015) with time series obtained from the bottom frame, we can 662 estimate the range of the SPM concentration variation resulting from the cross-shore oscillation. From the ship-based observations (van der Hout et al., 2015) a cross-shore SPM concentration 663 gradient 0.035 ± 0.015 mgl<sup>-1</sup> per meter at 1.2 mab is estimated (Figure 14). The trend of the vertical 664 665 profiles in Figure 14 suggests that the near-bed cross-shore gradient at 0.3 mab may be even 666 stronger. To estimate the cross-shore tidal water displacement a harmonic analysis was performed 667 on the velocity data from the bottom frame, followed by a primary and secondary axis 668 decomposition of the  $M_2$  component. The average amplitude of the minor axis appeared 0.02 ms<sup>-1</sup>, the maximum is 0.09 ms<sup>-1</sup>. The resulting cross-shore tidal water displacement is 280 m for the 669 average amplitude and 1300 m for the maximum amplitude. Multiplied by the average cross-shore 670 671 SPM concentration gradient, the average magnitude of the tidal SPM concentration variation at 1.2 672 mab due to cross-shore advection would be 9.8 mgl<sup>-1</sup>, and the maximum 65 mgl<sup>-1</sup>. Thus the cross-673 shore advection may contribute substantially to the variation in SPM concentrations when compared 674 to the median tidal variation of 84 mgl<sup>-1</sup> at 1.4 mab but is on average only 12%. This enforces our 675 previous suggestion that on the tidal scale resuspension is dominant over advection. Summarizing, 676 with the alternating tidal current the space-varying SPM field is moving up and down (resuspension -677 settling) while moving back and forth (advection) along the study site and possibly changing the 678 aggregation form of the flocs.

679



680

Figure 14. The vertical profiles of the horizontal SPM concentration gradient at the frame location. No data closer to the seabed is available than 1.2 mab. The SPM concentrations were measured during four surveys along the Egmond transect in 2010 and 2011 (Van der Hout et al., 2015), the dates of the visit are indicated in the legend. Positive SPM concentration gradients indicate an increase in concentration towards the shore; negative values indicate a decrease in SPM
concentrations towards the shore.

#### 687 **4.5.** *Rouse profiles*

688 A logarithmic Rouse profile fitted to 84% of the SPM profiles, but not to the CHL-a profiles. The 689 disparity in the quality of the fits of the CHL-a and SPM profiles suggests that although CHL-a is 690 aggregated in flocs of SPM, the aggregation is not uniform in the near-bed layer. A dependence of  $R_n$ on Ca is observed at our study site, similar to other studies (Orton and Kineke, 2001), which is a 691 692 violation of one assumptions for the Rouse profile. Apparently this underlying assumption is not a 693 strict requirement. Care should be taking with applying the Rouse profile at low SPM concentrations 694 and more specifically near slack tides, where most poorly-fitting profiles were found. Whereas Orton 695 and Kineke (2001) excluded a small period around slack (range of 0.01 ms<sup>-1</sup>) from their analysis, we 696 found that rejection occurred predominantly at the accelerating current after slack and up to a 697 velocity of 50 % of the maximum current. The sensitivity of the profile fit at low concentrations is 698 possibly due to the slow response of the low concentration particles to quickly changing conditions 699 (Orton and Kineke, 2001).

700 The calculated Rouse numbers showed that in 90 % of the time the SPM profile is a mixed 701 suspension, suggesting that SPM concentrations at the surface are in the same order of magnitude as 702 the near-bed concentrations. This implies that both the near-bed layer and the surface layer should 703 be taken into account when considering transport estimates of SPM in further research. The 704 contribution of the 10% bed load transport to the total transport should not be underestimated, as 705 these profiles with a high  $R_n$  also have the highest SPM concentrations ( $R_n \sim C_a$ ). Although 84% of the 706 SPM concentrations profiles could be fitted to the Rouse profile, extrapolating the near-bed data up 707 to the surface with the Rouse method should be done with care in this study area as the observed 708 periodic vertical stratification (Rijnsburger et al., 2016; Van der Hout et al., 2015) may temporarily 709 inhibit SPM resuspension up to the surface (Jones et al., 1998; Joordens et al., 2001).

# 710 *5. Conclusion*

From the analysis and discussion of the various variables and time scales we can summarize that in the TMZ a fluffy layer exists close to the seabed which consists of accumulated SPM and CHL-a rich material. The SPM consists mainly of fine silt with a high organic content, and appears dominantly in aggregated form. For significant wave heights larger than 0.8 m the resuspension by waves dominates the near-bed SPM and CHL-a concentration variation over the tidal resuspension. During storms the seabed can become depleted, after which the high concentrated fluffy layer forms again by the settlement of the temporarily redistributed particles. On the regular tidal scale a part of the SPM and CHL-a remains in the resuspension - settling mode, and is at the same time advected back and forth in both the alongshore and cross-shore direction. The seasonal CHL-a concentration is dominated by the spring bloom, but shows also resuspension by waves. The seasonal SPM concentration is primarily driven by the seasonal wave climate; high in autumn and winter, low in spring and summer. A secondary actor forms the seasonal activity of organic matter and biota, which have an opposing seasonal cycle to the SPM.

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